Blunt V-notch Brittle Fracture: An Improved Finite Fracture Mechanics Approach

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Abstract. The coupled Finite Fracture Mechanics (FFM) criterion is applied to investigate brittle fracture in rounded V-notched samples under mode I loading. The approach is based on the contemporaneous fulfilment of a stress requirement and the energy balance, the latter being implemented on the basis of a recently proposed analytical expression for the stress intensity factor. Results are presented in terms of the critical crack advance and the apparent generalized fracture toughness, i.e. the unknowns related to the system of two equations describing the FFM criterion. A validation of the theory is performed by employing varying root radius notched, as-quenched, AISI 4340 steel specimens fracture results.

Introduction

Several approaches have been successfully proposed to investigate the mode I brittle fracture of structural components containing rounded V-notches [1-7]. The common idea these criteria are based on is that failure takes place when the (punctual or averaged) stress or an energy-related quantity at a finite distance from the notch tip reaches a critical value: such a distance results to be a material constant.

The critical distance becomes a structural parameter if the stress and energy conditions are coupled one to each other, as it happens for both the Leguillon’s criterion and the FFM criterion [8-12]: whereas the former involves a punctual stress requirement, the latter is of an average type. The coupled criteria were introduced to remove some inconsistencies related to approaches based either on stress or on energy requirements [13].

Indeed, the stress condition can be easily implemented, since an accurate expression for the stress field along the blunt notch bisector was derived [14]. This allows to develop stress-based criteria [2,3,6] in a quasi-analytical way. On the other hand, as it regards the energy balance, the relationships proposed for the stress intensity factor (SIF) related to a crack stemming from a blunt V-notch tip (providing the crack driving force) are usually approximated. Although closed-form solutions can be obtained, the results are affected by some notch amplitude depending errors, the entity of which has to be discussed case-by-case [11]. An alternative way may consist in proceeding numerically, via a specific finite element analysis [1,9,10]. Of course, this implies a more laborious computation, since it is necessary to carry out a different simulation for each data set. Similar considerations hold for the finite-volume-energy based approach [4].

As it concerns approximating functions for the SIF of a crack stemming from a blunt V-notch, let us mention Lukas’s formula [15] providing an approximating function for the SIF of a crack emanating from an elliptical hole. A U-notch can be considered as a limit case of an elliptical hole (see also [16]), when the minor axis to major axis length ratio tends to zero; upon observing that no significant deviations are generally expected for opening angles \(\omega\) \((0° \leq \omega \leq 90°)\), Lukas’ formula can be implemented [5-12], despite its range of validity is very limited. An analytical relationship
for the SIF related to a rounded V-notch was further proposed to fulfil the asymptotic limits of very short/long cracks [11]. It was numerically verified for $\omega = 90^\circ$, $120^\circ$, and $150^\circ$, showing the maximum deviation (nearly 7%) at $\omega = 90^\circ$. Since the energy available for a crack length increment is proportional to the integral of the squared SIF, the global error is expected to increase. Even more recently, the relationship was improved by adding a notch amplitude parameter $m$, the value of which was fitted by means of ad hoc numerical simulations [17]: in this case, the error on the SIF was estimated to be lower than 1% over the whole range $0^\circ \leq \omega \leq 180^\circ$, which can be considered very satisfying for practical applications.

The relationship provided in [17] is hereby exploited to enhance the FFM approach. The criterion, introduced in Section 2, consists in a system of two equations (one for the stress requirement and the other for the energy balance, respectively) in two unknowns: the apparent generalized fracture toughness (i.e., the failure load) and the critical crack advance. As already hinted, the latter, which coincides with the critical distance, results to be a structural parameter, controlled by the notch amplitude and the root radius, as well as by the material properties. Dimensionless results are presented in Section 3, together with the comparison with experimental data [18-20].

It has to be noted, that FFM has been already applied to the specific case of U-notches [12] by implementing Lukas’ formula and to the general framework of structures containing blunt V-notches [11]. The novelties of the present work with respect to the past ones are substantially two: (i) the FFM implementation via the relationship proposed in [17], what allows to improve the prediction accuracy, especially on samples with smaller notch amplitudes, as well as to develop a unified FFM approach in the framework of blunt V-notches. Indeed, since the crack driving function can not be integrated explicitly, numerical quadrature formulae must be implemented; (ii) the comparison of FFM results with various experimental data sets, covering the whole $0^\circ \leq \omega \leq 150^\circ$ range.

### The Finite Fracture Mechanics Criterion

The FFM criterion is based on the hypothesis of a finite crack advancement $l$ and assumes the contemporaneous fulfilment of two conditions [21-22]. By referring to the coordinate system displayed in Fig. 1a, the criterion can be expressed as:

$$
\begin{align*}
\int_0^l \sigma_y(x) \, dx & \geq \sigma_u \, l \\
\int_0^l G(c) \, dc & \geq G_c \, l
\end{align*}
$$

(1)

The former condition requires that the average stress $\sigma_y(x)$ upon the crack advancement $l$ is higher than the material tensile strength $\sigma_u$. The latter one ensures that the energy available for a crack length increment $l$ (involving the integration of the crack driving force $G$ over such a length) is higher than the energy necessary to create the new fracture surfaces, $G_c \, l$. By means of Irwin’s formula, conditions (1) can be recast, at the critical mode I point, in the following form:

$$
\begin{align*}
\int_0^l \sigma_y(x) \, dx & = \sigma_u \, l_c \\
\int_0^l K_{Ic}^2(c) \, dc & = K_{Ic}^2 \, l_c
\end{align*}
$$

(2)

$K_I(c)$ and $K_{Ic}$ being the SIFs related to a crack of length $c$ stemming from the notch root (Fig.1b) and the fracture toughness, respectively. A system of two equations in two unknowns is obtained, with the unknowns being the critical crack advancement and the failure load (implicitly embedded in the stress field and SIF functions). Note that, upon passing from (1) to (2), the hypothesis of positive geometries (monotonic increasing $K_I(c)$ functions) has been implicitly assumed, allowing one to substitute the inequalities with the corresponding equalities.
Fig. 1 (a) Blunt V-notch with Cartesian and polar coordinate systems; (b) crack of length $c$ stemming from the notch tip.

In Fig. 1, the distance between the notch tip and the origin of the polar coordinates system is indicated as $r_0$:

$$r_0 = \frac{2\pi - \omega}{\pi - \omega} \rho$$

(3)

where $\rho$ is the notch root radius.

For sharp cracks, the origin coincides with the notch root end whatever the value of $\omega$. As the crack becomes blunter, $r_0$, which is also the distance between the Cartesian and polar coordinate systems origins, increases for a given value of $\omega$. Conversely, it also increases as the notch amplitude increases for a given value of notch root radius.

The stress field. By assuming that the notch tip radius $\rho$ is sufficiently small in respect to the notch depth $a$, the stress field along the notch bisector could be expressed in polar coordinates as in [14]:

$$\sigma_\theta(r, 0) = \frac{K_i^V}{2\pi r^{1.5}} \left[ 1 + \left( \frac{r_0}{r} \right)^{1-\mu} \eta_\theta(0) \right]$$

(4)

where $K_i^V$ is the generalized stress intensity factor referring to a V-notch of depth $a$. The eigenvalues $\lambda$ and $\mu$, as well as the function $\eta_\theta(0)$, depend on the notch amplitude $\omega$. These latter three parameters are reported in Fig. 2.

The SIF function. Let us consider a crack of length $c$ stemming from a blunted V-notch root (Fig. 1b). As far as the notch depth $a$ is sufficiently large with respect to $c$, the following SIF function was proposed [17],

$$K_i(c) = \frac{K_i^V \beta c^{1/2}}{1 \left( \frac{\beta}{\psi} \right)^{1-\mu} \left( \frac{1}{r_0} \right)^{1-\mu}}$$

(5)

The notch angle amplitude functions $\beta$ and $m$ are drawn in Fig. 2 as a function of $\omega, \psi$ is equal to $1.12\sqrt{\pi} (1+\eta_\theta(0)) (2\pi)^{1-\mu}$. Note that the parameter $m$ was fitted to improve the predictions obtained by carrying out an ad hoc finite element analysis [17]: Eq. (5) fulfils the asymptotic limits of very short and very long (but still small in respect to the notch depth) cracks, providing errors below 1%, for $0 \leq c/\rho \leq 10$ and over the range $0^\circ \leq \omega \leq 180^\circ$. 

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**Fig. 1** (a) Blunt V-notch with Cartesian and polar coordinate systems; (b) crack of length $c$ stemming from the notch tip.
FFM implementation

By hypothesizing that failure takes place when the generalized SIF, \( K'_{V} \), reaches its critical conditions \( K'_{V}^{c} \), as expected for brittle structural behaviour, FFM can be implemented by inserting Eqs. (4) and (5) into system (2) and integrating. Note that \( K'_{V}^{c} \) represents the apparent generalized fracture toughness (i.e. the generalized fracture toughness measured as if the V-notch was sharp): it depends on the root radius, differently from the generalized SIF, \( K'_{I} \). Some analytical manipulations yield:

\[
\begin{align*}
\frac{K'_{V}^{c}}{\sigma_{u}^{k-2}} &= f(\bar{T}_{c}), \\
\frac{K'_{V}^{c}}{K'_{I}^{c}} &= h(\bar{T}_{c}),
\end{align*}
\]

where,

\[
\begin{align*}
f(\bar{T}_{c}) &= \frac{(2\pi)^{1-\lambda} \bar{T}_{c}}{\left[\left(\bar{T}_{c}+1\right)^{\lambda} - 1\right] / \lambda + \eta_{0}(0) \left[\left(\bar{T}_{c}+1\right)^{\mu} - 1\right] / \mu}, \\
h(\bar{T}_{c}) &= \frac{\bar{T}_{c}}{\beta c^{2\lambda - 2\eta - 1}} \int_{0}^{1} \frac{1}{1 + \left[\frac{\beta}{\psi}\left(\frac{1}{c}\right)^{\mu - 1} \left(\frac{1}{\bar{T}_{c}}\right)^{\mu} \right]^{m}} d\bar{\sigma},
\end{align*}
\]

\( \bar{T}_{c} = l_{c} / r_{0} \) and \( \bar{\sigma} = c / r_{0} \).

Note that the presence of the parameter \( m \) in Eq. (7b) does not allow to get a closed form solution for function \( h \): the integral will be thus evaluated numerically according to a recursive adaptive Simpson quadrature formula. By equaling both expressions in (6) with respect to \( K'_{V}^{c} \), the problem can be recast in the following form:

![Figure 2 Notch angle amplitude dependent parameters used in the present analysis.](image-url)
\[
\begin{align*}
\frac{K_{lc}^{V,\rho}}{\sigma_u l_0^{1-\lambda}} &= f(\bar{l}_c), \\
r_0 &= \frac{h(\bar{l}_c)}{f^2(\bar{l}_c)}, \\
l_{ch} &= \frac{2}{\lambda \beta^2 (2\pi)^{2(1-\lambda)} l_{ch}}, \\
\end{align*}
\]  

(8)

where \( l_{ch} = (K_{lc} / \sigma_u)^2 \) is the Irwin’s length. For a given material, notch amplitude \( \omega \) and notch root radius \( \rho \), the value of the critical crack advancement \( \bar{l}_c \) can be derived from the second equation in (8). This value must then be inserted into the former equation to obtain the apparent generalized fracture toughness, \( K_{lc}^{V,\rho} \).

FFM results are presented in Figs. 3 and 4, in dimensionless form as compared to the sharp V-notch case, according to which:

\[
l_c^V = \frac{2}{\lambda \beta^2 (2\pi)^{2(1-\lambda)} l_{ch}},
\]

(9a)

\[
K_{lc}^V = \lambda \left[ \frac{2(2\pi)^{2\lambda-1}}{\beta^2} \right]^{1-\lambda} l_{ch}^{1-\lambda} \sigma_u.
\]

(9b)

Fig. 3 Dimensionless plot of the critical crack advancement vs notch root radius according to FFM.

Fig. 4 Dimensionless plot of the apparent generalized fracture toughness vs. notch root radius according to FFM.
In Fig. 5, the apparent generalized fracture toughness is reported according to the present FFM approach and to the one proposed in [11], obtained simply by setting $m=1$ in Eq. (5). Note that the maximum deviations (and consequently the most significant improvements) are observed for lower root radii and lower notch angle amplitudes $\omega$, as expected since $m$ decreases as $\omega$ increases.

![Graph showing different FFM predictions for apparent generalized fracture toughness](image)

Fig. 5 Different FFM predictions for the dimensionless apparent generalized fracture toughness: the thin lines describes the results obtained in [11], the thick lines are related to the present improved approach. From top to bottom, curves refer to $\omega=0^\circ$, $30^\circ$, $60^\circ$, $90^\circ$, $120^\circ$ and $150^\circ$.

**Comparison with experimental data.** Let us now compare FFM predictions with experimental data available in the literature: different materials, root radii and notch amplitudes are taken into account.

The firstly considered data set refers to three point bending tests on single edge U-notched specimens made of alumina [20]. The material properties are $K_c = 3.81 \text{MPa}\sqrt{\text{m}}$ and $\sigma_u = 290 \text{MPa}$, corresponding to $l_{ch} \approx 0.173 \text{mm}$. FFM predictions, in terms of the apparent fracture toughness $K_{cU} = K_c^{\omega=0,\rho}$ are plotted in Fig. 6, showing a very good matching. The calibration of the parameter $m$ in Eq. (5) allows to get a more flat behaviour in the neighbourhood of $\rho = 0$ with respect to the FFM approach put forward in [11]. Note that the presence of a minimum was highlighted in [7] by implementing some criteria based on a critical distance.

The second data set taken into account is related to Charpy V notch specimens fabricated with as quenched AISI 4340 steel [18]: they were heat treated by a 1200°C- 870°C step quenching procedure, yielding a martensitic, 250 \text{mum} prior austenitic grain size microstructure, and tested under three-point bending. Although the V-notch angle related to Charpy samples is $45^\circ$, it was considered as a crack ($\omega = 0^\circ$) in [18-19]. The measured AISI 4340 steel properties are: $K_c = 73.9 \text{MPa}\sqrt{\text{m}}$ and $\sigma_u = 1980 \text{MPa}$, resulting in $l_{ch} \approx 1.39 \text{mm}$. The material is of course less brittle than common ceramic materials, usually tested to investigate blunt V-notch brittle fracture. Indeed, it has been shown that some criteria, such as the point-stress approach, work rather well even for ductile metallic materials, when the final rupture is preceded by large-scale plastic deformations [23]. FFM results for as quenched AISI 4340 large prior austenitic grain size steel are shown in Fig. 7.
Conclusions

The novel expression for the SIF related to a crack stemming from a blunted V-notch proposed in [17] allows us to improve FFM predictions on mode I brittle fracture. The analysis covers a wide range of notch radii and two different types of materials. Results are in good agreement with experimental data obtained by the use of a truly brittle material such as alumina, whereas large notch root radii samples fracture results, obtained with a martensitic, large prior austenitic grain size.
steel, exhibit a deviation from the model, thus imposing further analysis in the case of less brittle materials as well as large notch root radii.

As a future development, the next step will consist in investigating the effects of the root radius on mixed-mode brittle fracture through FFM: note that some criteria have already been successfully proposed [24]. Some considerations have been already pointed out also through FFM [25], although the direct implementation of the approach in this framework is still missing.

References
