Energy harvesting from wind by a piezoelectric harvester

J.X. Tao a, N.V. Viet b, A. Carpinteri c, Q. Wang d, ⇑

a City University of Hong Kong Shenzhen Research Institute Building, Shenzhen Hi-Tech Industrial Park, Nanshan District, Shenzhen, China
b Department of Mechanical Engineering, Khalifa University, P.O. Box 127788, Abu Dhabi, United Arab Emirates
c Dept. of Structural, Geotechnical and Building Engineering, Politecnico di Torino, Corso Duca degli Abruzzi, Italy
d Department of Architecture and Civil Engineering, City University of Hong Kong, Kowloon, Hong Kong, China

Abstract

A novel wind power harvester using piezoelectric effect is developed. This engineering structure device employs a wind turbine that extracts kinetic energy from the wind and converts it into the rotational motion of a shaft. The rotational shaft is connected to a Scotch yoke mechanism that is used for converting the rotational motion into linear vibrations of two piezoelectricity-levers through springs. A mathematical model is developed to calculate the root mean square value of the generated electric power. The influences of some practical considerations, such as the rotational speed of the wind turbine and the stiffness of the springs, on the root mean square of the generated power are discussed. The research results show that a power up to 150 W can be harvested for a piezoelectric wind turbine with a radius of blades of 1 m at the wind speed of 7.2 m/s, and the designed angular velocity of 50 rad/s.

Keywords:
Energy harvesting
Wind turbine
Piezoelectric harvester

1. Introduction

The increasingly consumption of traditional energies has caused serious energy crisis and environmental pollution on a global scale. Wind energy has been regarded as one of the most important renewable and green energy sources to solve the above problems [1]. Wind turbines have been used on one of the traditional ways of transferring kinetic energy from wind based on the principle of electromagnetic induction [2,3]. The power coefficient of the wind turbines with output of megawatt has been reported up to 40–45% at the optimal tip speed ratio of 5–7 [4]. However, these large-size turbine generators are always complex, costly, and normally induce a cogging torque which restricts the cut-in speed. For small wind turbines located in areas of unfavorable wind, they are of little practical use except in rare situations [5,6]. As a result, using unconventional approaches such as vortex induced vibration, galloping, flutter, and buffeting to harvest the wind energy have been recently proposed [7–15]. To provide a more efficient energy transfer process, piezoelectric technology is the most prominent candidate for converting mechanical energy to electricity with highest efficiencies and lowest costs owing to their high energy generation density, high voltage generation capability and simple configurations and economic benefits [16–18].

Early works about a piezoelectric windmill were carried in literatures [9,10]. The proposed devices had ten piezoelectric bimorphs arranged along the circumference of a horizontal-axis wind turbine rotor shaft in the cantilever beam form. The oscillating torque to vibrate the bimorphs was generated using the camshaft gear mechanism. A power of 7.5 mW at the wind speed of 10 mph was measured across a matching load of 6.7 kΩ. Authors also addressed some drawbacks of this device and gave an optimized structure made of only plastic parts in the later work [11]. Sirohi et al. [12] developed a piezoelectric energy harvesting device based on a galloping cantilever beam. The harvested wind energy is transferred to a galloping beam which has a rigid tip body with a D-shaped cross section. Piezoelectric sheets were bonded on the top and bottom surface of the beam. During galloping, vibrational motions are induced due to aerodynamic forces on the D-section, which is converted into electrical energy by the piezoelectric (PZT) sheets. Their experimental and analytical investigations of dynamic response and power output have shown that a maximum output power of 1.14 mW was measured at a wind velocity of 10.5 mph on a prototype device of length 235 mm and width 25 mm. Rezaei-Hosseinabadi et al. [13] presented a topology for piezoelectric energy harvesting made of a lift-based wind turbine and a piezoelectric beam with contactless vibration mechanism. The research results showed that a power density of 2 mW/cm³ at 3.8 V at the wind speeds above 0.9 m/s can be achieved. Kishore et al. [14] designed an ultra-low start-up speed windmill made of a 72 mm diameter horizontal axis wind turbine rotor with 12 alter-
nating polarity magnets around its periphery and a 60 mm × 20 mm × 0.7 mm piezoelectric bimorph element having a magnet at its tip. This wind turbine was found to produce a peak electric power of 450 μW at the rated wind speed of 4.2 mph. Wu et al. [15] developed an effective and compact wind energy cantilever harvester subjected to a cross wind. Sufficient electrical energy output as high as 2 W was realized by tuning the resonant frequency of the harvester with a proof mass on the tip of the cantilever.

The aforementioned piezoelectric harvester devices have been mainly developed for energy supplies of wireless sensors. Therefore, most of them are small in sizes and the energy outputs are always in the scale of mW. Recent researches have shown that piezoelectric harvesters can harvest power range from several watts to hundreds of megawatts [19–23]. Xie et al. [19,22] developed a novel piezoelectric technology of energy harvesting from high-rise buildings and found that a power up to 432.2 MW can be realized. Viet et al. [23] proposed a floating energy harvester (FEH) using piezoelectric effect to harvest energy from water waves. Based on their simulated results, the root mean square (RMS) of 103 W can be achieved when the wave amplitude is 2 m. Hence, it is imperative to develop harvesters with high power using piezoelectric effect from wind.

In this research, a wind energy harvester using piezoelectric technique is developed. This proposed piezoelectric wind turbine incorporates the advantages of the conventional wind turbines and the piezoelectric harvesters. Wind turbine rotor blades with the shape optimized by aerodynamics and connected to a horizontal shaft are employed as the driving device. The wind-induced rotational motion of the shaft is converted into translation motion of a slotted rod through a Scotch yoke mechanism. Both sides of the slotted rod are linked with a spring used for transferring the linear motion into vibrations of two piezoelectric levers. The Scotch yoke, shown in Fig. 2, is a reciprocating motion mechanism, converting a rotational motion into a linear motion [24]. A slotted rod is used to make sure the motion is only in the direction perpendicular to the axis of the shaft by a cylindrical slider attached to a wheel. When the wheel rotates at an angular velocity, \( \omega \), the end points of the slotted rod are displaced from their initial position by an amount \( z_t \) (in time \( t \)) given by \( z_t = Y \sin(\omega t) \), where \( Y \) is the amplitude, i.e. the distance between axes of the cylindrical slider and the shaft.

The piezoelectricity-lever device, shown in Fig. 3(a), consists of a lever with a long moment arm \( L_2 \) and a short moment arm \( L_1 \), a fixed-hinge for restricting linear displacements of the lever, and a piezoelectric bar with a Young’s modulus, width, length, and height of \( E_p \), \( a \), \( b \), and \( h \), respectively. Since the piezoelectricity-lever devices are connected to the slotted rod by two springs, the harmonic motion is then converted into a spring force \( F(t) \). The force location is at the point C and magnified \( n \) times at the point A on the piezoelectric bar by the lever mechanism, where \( n \) is the ratio of the length of the moment arm to that of short moment arm, \( n = L_2/L_1 \). Consequently, the electric power is generated by the piezoelectric-lever design.

In order to estimate the spring force \( F(t) \), one of the piezoelectric lever is used as an example and is simplified as a damped single-degree-of-freedom (spring-mass) system, shown in Fig. 3(b) [23,25]. The equivalent mass \( m_e \), spring stiffness \( k_e \), and damping coefficient \( c \) can be derived from the material properties and dimensions of the lever and piezoelectric bar.

In the proposed design, the piezoelectric bar is firmly bonded on the lever at the region of point A. Due to the high tensile/compression stiffness of the piezoelectric bar in the force direction, the lever is supposed to be fixed throughout points A and B and equivalently transformed into a cantilever beam with a length of \( L_2 \). Hence, the equivalent mass \( m_e \) can be calculated as \( m_e = \rho AL_2/n \), where \( \rho \) and \( A \) are material density and cross-section area, respectively. In this work, the lever is made of steel, and its cross-section is rectangular in shape with a width and height of \( S_1 \) and \( S_2 \), respectively. The dimensions of the lever face are selected as \( s_1 = 0.015 \) m and \( s_2 = 0.004 \) m, unless otherwise noted.

Since the applied force at point C induces an elastic deflection of the cantilever and axial deformation of the piezoelectric bar, the equivalent spring stiffness, \( k_e \) in Fig. 3(b) can be yielded as [26]:

![Fig. 1. Schematic diagram of a horizontal piezoelectric wind turbine: (a) an overview of the device, and (b) an internal structure of the device.](image-url)
function of the single-degree-of-freedom system.

\[
k_e = k_h/(k_0 + k_p), \quad k_i = E_0 s_0^3/(4L_1 n)^2, \quad k_p = E_p ab/(\pi^2 h)
\]  

(1)

where \( k_0, E_0, k_p, E_p \) are spring constant and Young’s modulus of the cantilever and piezoelectric bar, respectively.

The equivalent damping coefficient, shown in Fig. 3(b), can be expressed as a sum of the mechanical and electrical damping coefficients:

\[
c = c_i + c_e
\]  

(2)

The mechanical damping force of the cantilever beam is triggered due to the internal structure damping force [27] and the friction between the viscous air and the cantilever beam in the vibration process, and hence is given as:

\[
c_i = 2\xi \sqrt{k_m l}
\]  

(3)

where \( \xi \) is the damping ratio.

Besides the mechanical damping, there is an equivalent electric resistance when the work done by the spring force is converted into electricity [28]. The electrical damping coefficient, \( c_e \), can be derived as [20]:

\[
c_e = \pi^2 d_{33}^2 k_e^2/(\pi^2 c_a f)
\]  

(4)

where \( d_{33} \) is the piezoelectric constant in the polling direction; \( c_a \) is the electric capacity of the piezoelectric bar; \( f \) is the natural vibration frequency of the single-degree-of-freedom system.

Once the equivalent mass \( m_e \), spring stiffness \( k_e \), and damping coefficient \( c_e \) are obtained, we can get the equation of motion of the equivalent mass \( m_e \) using Newton’s second law:

\[
m_e \ddot{z}_e + c_e \dot{z}_e + k_e z_e = (z_t - z_e) k_1
\]  

(5)

where \( z_t \) and \( z_e \) are the displacements of the slotted rod and the mass \( m_e \), respectively. As is discussed above, horizontal motion of the slotted rod is expressed as \( z_t = Y \sin \omega t \). Eq. (5) is then transformed as:

\[
m_e \ddot{z}_e + \ddot{z}_t + (k_e + k_1) z_e = z_1 k_1 = k_1 Y \sin \omega t
\]  

(6)

Hence, the model can be considered as a damped single-degree-of-freedom system subjected to an equivalent harmonic force:

\[
m_e \ddot{z}_e + \ddot{z}_t + k_e z_e = F \sin \omega t
\]  

(7)

where \( k = k_e + k_1 \) and \( F = k_1 Y \). We can easily obtain the displacement of equivalent mass \( m_e \) [29]:

\[
z_e(t) = X \sin(\omega t - \phi)
\]  

(8)

where \( X \) is a constant that denotes the maximum amplitude of \( z_e(t) \); \( \phi \) is the phase angle:

\[
X = \frac{F}{\sqrt{(k - m_0 \omega^2)^2 + (\omega \phi)^2}}
\]  

\[
\tan \phi = \frac{\omega \phi}{k - m_0 \omega^2}
\]  

(9)

Consequently, the relative displacement, \( z(t) \), of the spring-mass system is:

\[
z(t) = z_t(t) - z_e(t)
\]  

(10)

The generated electric power on each piezoelectric bar can be expressed below:

\[
P_e(t) = I(t) \cdot V(t)
\]  

(14)
Since two piezoelectricity-lever devices are included in this device, the total generated electric power should be:

\[ P_s(t) = 2P_e(t) \]  \hspace{1cm} (15)

Finally, the expression of RMS of the electric power is provided. The RMS of the generated electric power from time 0 to \( t \) can be given as:

\[ P_{s}^{RMS} = 2 \sqrt{\frac{1}{T} \int_{0}^{T} P_s^2 dt} \]  \hspace{1cm} (16)

After the mathematic model is developed, we can estimate the generated charge and voltage from the piezoelectric bar as well as the RMS of the electric power. In this work, the material properties and dimensions of the piezoelectricity (PZT4)-lever device are provided in Table 1, unless otherwise noted. To find an efficient design of the piezoelectricity-lever and verify the effectiveness of the harvesting system, the aerodynamic model of wind turbine are adopted here. The mechanical power of the turbine, \( P \), is given by [6]:

\[ P = \frac{1}{2} \rho \pi R^2 v_w^3 C_p(\lambda) \]  \hspace{1cm} (17)

where \( \rho \) is the air density, which equals 1.225 kg/m\(^3\) at a temperature of 15 °C at sea level; \( R \) is the turbine radius; \( v_w \) is the wind speed. The factor \( C_p(\lambda) \) is the power coefficient of the turbine, which is determined by the Tip-Speed Ratio \( \lambda \). The tip-speed ratio \( \lambda \) is calculated as the ratio between the speed of the tip of a blade and the wind speed \( v_w \):

\[ \lambda = \frac{R \omega}{v_w} \]  \hspace{1cm} (18)

The relation between the power coefficient \( C_p \) and the tip-speed ratio \( \lambda \) is determined by the shape of the blade. An empiric \( C_p(\lambda) \) relations exist in literature [30] is used here:

\[ C_p(\lambda) = 0.73 \left( \frac{151}{\lambda} - 13.65 \right) e^{\left( -\frac{2}{\lambda-0.055} \right)} \]  \hspace{1cm} (19)

In the traditional horizontal-axis wind turbines, blades rotates only when the aerodynamic torque, \( T \), exceeds the combined resistive torque of the drive train and generator, \( T_g \). The starting rotating wind speed, when \( T = T_g \), is called the cut-in speed. Wood [31] used the standard blade element theory to calculate the cut-in wind speed and provided the following equation:

\[ U_c = \left( \frac{2T_g}{NpR^2d_{33}} \right)^{1/2} \]  \hspace{1cm} (20)

where \( U_c \) is the cut-in speed; \( T_g \) is the resistive torque and can be expressed as \( T_{max} = k_1Y^2 \) in this model; \( N \) is the number of blades; \( \rho \) is the air density; \( R \) is the blade tip radius; and \( d_{33} \) is the chord-pitch integral determined by the shape of the blades.

From Eq. (20), we can find that the cut-in speed increases with the increasing resistive torque for the same blades. Therefore, for wind turbines located in areas of unfavorable wind, the resistive torque should be very small. In addition, increasing the number and the radius of the blades lead to a low cut-in speed for the same resistive torque.

### 3. Results and discussions

In this section, based on the above developed numerical model, we investigate the effectiveness of the piezoelectric wind turbine. Particularly, this work studies the effect of the spring stiffness, the amplitude, the angular velocity of the wind turbine, the ratio of moment arms, and the width (length) of the piezoelectric bar on the generated power. By using the aerodynamic model and empiric relations between the power coefficient \( C_p \) and the tip-speed ratio \( \lambda \), the extracted mechanical power of a three-bladed wind turbine with radius of 1 m can be calculated, and the simulated results are compared with the RMS of the generated electric power.

The effect of spring stiffness \( k_1 \) and the amplitude \( Y \) on the RMS are first investigated because they determine the performance of the piezoelectric wind turbine. As mentioned before, for small wind turbine located in areas of unfavorable wind, the resistive torque of the electric generator should be very small. The field test results in the work of Clausen et al. [5] showed that for a designed two-bladed turbine with the radius of 2.5 m, the cut-in speed is less than 3 m/s when the resistive torque is 1 Nm. Wood [31] used a derived equation to calculate the cut-in speed and provided a value of 4.75 m/s for a three-bladed turbine with a length of 1 m. In this model, several small values ranging from 0.4 Nm to 1.2 Nm, of resistive torque, \( k_1Y^2 \), are used. Fig. 4 gives the effect of spring stiffness \( k_1 \) on the RMS of the generated electric power when the angular velocity, ratio of moment arms, and the width of the piezoelectric bar are 50 rad/s, 20, and 0.015 m, respectively. It is revealed from Fig. 4 that, when \( k_1Y^2 \) is set as a constant, the RMS increases nonlinearly with a rise of the spring stiffness. When the stiffness changes from 5000 N/m to 40,000 N/m, the RMS increases from 15 W and 30 W to 60 W and 180 W corresponding to the resistive torque of 0.4 Nm and 1.2 Nm, respectively. This observation is interpreted by the fact that the force applied to the piezoelectric bar is proportional to \( k_1Y \) according to Eq. (11). When the resistive torque \( k_1Y^2 \) keeps a constant, a growth in the stiffness of springs will lead to a decline in the value of \( k_1Y \), shown in Fig. 5, and correspondingly an enhancement in the force applied to the piezoelectric bar. The finding indicates that increasing the stiffness of the spring and decreasing the amplitude to keep a constant resistive torque can get a higher energy harvesting efficiency. In addition, increasing the amplitude \( Y \) and keeping a constant stiffness of the spring can also enhance the applied force \( k_1Y \). But it should be noted here, the resistive torque will increase accordingly, which will make the turbine less effective at low speed wind regions.

The effect of the angular velocity of the wind turbine on the RMS is shown in Fig. 6. The stiffness of the spring, the amplitude, and the ratio of moment arms are set as 40,000 N/m, 0.005 m, and 20, respectively. It is seen clearly that the RMS increases nonlinearly from 0 rad/s to 252 W with a rise in the angular velocity from 0 rad/s to 100 rad/s. This observation is mainly because that an increase in the angular velocity leads to a decrease in the period of the equivalent harmonic force, and hence causes an increase in the kinetic energy of the harvesting system and the generated electric power. Actually, the optimal angular velocity can be determined by Eqs. (17)–(19). In the real design, a large power coeffi-

| Material properties and dimensions of the piezoelectricity (PZT4)-lever device. |

<table>
<thead>
<tr>
<th>Piezoelectricity (PZT4)-lever device</th>
<th>( a ) (m)</th>
<th>( b ) (m)</th>
<th>( h ) (m)</th>
<th>( l_1 ) (m)</th>
<th>( \zeta )</th>
<th>( d_{33} ) (C/N)</th>
<th>( E_p ) (N/m²)</th>
<th>( E_l ) (N/m²)</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.015–0.025</td>
<td>0.015–0.025</td>
<td>0.1</td>
<td>0.02</td>
<td>0.0017</td>
<td>3e–10</td>
<td>7.5e10</td>
<td>2.1e11</td>
<td></td>
</tr>
</tbody>
</table>

\( C_v \) (nF) 0.375 for the piezoelectric patch with the geometry of 0.01 m, 0.01 m, 0.0001 m.
As has been noted, wind turbine starts to rotate only when the wind speed is larger than the cut-in speed. For a given resistive torque, the cut-in speed can be estimated based on Eq. (20). Using the chord-pitch integral in Ref. [31], the calculated cut-in speed for a three-bladed turbine with a radius of 1 m is 7.5 m/s when the resistive torque is 1 Nm. This value over-estimates the actual cut-in speed by a factor of two according to the conclusion of Wood [31], which means the actual cut-in speed is nearly 4 m/s. Therefore, the designed wind speed ranging from 5.6 m/s to 8.5 m/s is used to study the efficiency of the proposed device. The simulation results are depicted in Fig. 7.

The results indicate that, keeping the optimal tip-speed ratio, a wind turbine with a radius of 1 m can extract the maximum mechanical power ranging from 152 W to 553 W when the wind speed varies from 5.6 m/s to 8.5 m/s, and the its corresponding angular velocity changes from 39 rad/s to 60 rad/s. Comparing the RMS of electric power in Fig. 6 with mechanical power in Fig. 7, we find that for the same angular velocity, the extracted wind power of the turbine is always higher than the RMS of the harvested electric power. For example, the wind turbine achieves a power of 308 W when the designed wind speed and angular velocity are 7.2 m/s and 50 rad/s, respectively. However, the RMS of the electric power is only 150 W when \( \omega = 50 \) rad/s, which is nearly half of the power extracted by the wind turbine. In addition, this gap becomes larger with an increase in the angular velocity. This finding means the efficiency of the proposed device at large angular velocity is lower than that at small values. The main reason for the low efficiency is that the resistive torque in Fig. 6 is only 1 Nm, and the low resistive torque leads to a small force applied to the piezoelectric bar. Mechanical and electric damping existed in the piezoelectricity-lever also makes the present harvester less effective at large angular velocity. Therefore, in the real practice, wind turbine with a larger radius is recommended in order to get a larger starting torque. Besides, a gearbox used to amplify the angular velocity can be adopted for generating a high electric power.

The effect of the ratio moment arms of the lever on the RMS of the electric power generated by the piezoelectric harvester can be clearly seen in Eqs. (11) and (12). Since the applied force is proportional to the ratio moment arms of the lever, a larger ratio of the moment arms of the lever will lead to a higher power generated by the piezoelectric bar. Fig. 8 gives the displacement of the equivalent mass \( m_c \) when \( n = 20 \). Wherein, the angular velocity, the stiffness of the spring, and the amplitude are 50 rad/s, 40,000 N/m, and 0.005 m, respectively. Since the maximal displacement in Fig. 8 is only 0.00094 m, it is reasonable to use a large ratio \( n \) in the real design which will not break the lever either.

Fig. 9 illustrates the effect of the width and length of the piezoelectric bar on the RMS of electric power. The angular velocity, the stiffness of the spring, the amplitude, and the ratio moment arms are set to be 50 rad/s, 40,000 N/m, 0.005 m, and 20, respectively. It can be found that the RMS nonlinearly decreases with an increase in the width and length of the piezoelectric bar. The RMS decreases from 150 W to 74 W when the width and length of the piezoelectric bar change from 0.015 m to 0.025 m. The findings show that an increase in the width and length of the piezoelectric bar will lead to a rise in the electric capacity of \( c_{nr} \) and in turn a decrease in the electric voltage.
4. Conclusion

In this research, a practical piezoelectric wind turbine based on an engineering structure device is developed to harvest energy from wind. The Scotch yoke mechanism is adopted which is driven by rotating blades subjected to a force generated by wind. Hence, the rotational motion is converted into the linear motion of a slotted rod which is connected with two levers by springs. On the other side of the levers, there are two piezoelectric bars. The piezoelectric harvester transfers and amplifies the collected mechanical energy to electrical power. A mathematic model is developed to calculate the RMS of the generated electric power. The results show that, for a constant resistive torque, the RMS increases with an increase in the stiffness of the springs, angular velocity of the wind turbine, and the ratio of moment arms of the lever. Some experiential and theoretical formula of wind turbine are adopted to study the efficiency of the proposed device. The simulation results show that a power up to 150 W can be harvested for a piezoelectric wind turbine with a radius of blades of 1 m at the wind speed of 7.2 m/s, when the designed angular velocity is 50 rad/s. Based on the findings, it is possible to harvest a higher power by increasing the number and the radius of wind turbine rotor blades.

References